# Problem 12 Resonating Glasses 

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## Problem

When you take two glasses between your fingers, they sometimes emit a particular sound containing a frequency sweep. Investigate the phenomenon.

To investigate phenomena 1. Simplify Model
2. Analyze Data
3. Extend Considerations

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## Simplify as a problem of pendulums

Rigid body: For visual purposes, we draw the 'ball' at the center of mass.


Further reduce to single pendulum \& wall


## Relevant parameters

 for successive collisions.$\rightarrow \quad \boldsymbol{\theta}_{\mathrm{k}} \sim$ max angle before $\mathrm{k}^{\text {th }}$ collision
$\rightarrow \quad \boldsymbol{\omega}_{\mathrm{k}} \sim$ angular velocity immediately after the $(k-1)^{\text {th }}$ collision


## What does the system look like for energy?

|  | Initial Max Angle | Immediately <br> Before Collision | Immediately <br> After Collision | Next Max Angle |
| :--- | :---: | :---: | :---: | :---: |
| Kinetic Energy | 0 | $\mathrm{KE}_{\mathrm{k}-1}$ | $\mathrm{KE}_{\mathrm{k}}=\mathrm{c}^{*} \mathrm{KE}_{\mathrm{k}-1}$ | 0 |
| Potential Energy | $\mathrm{L}^{*} m g\left(1-\cos \left(\theta_{\mathrm{k}-1}\right)\right)$ | 0 | 0 | $L^{*} m g\left(1-\cos \left(\theta_{\mathrm{k}-1}\right)\right)$ |

Using the small angle approximation along with $\mathrm{W}=\triangle \mathrm{KE}$

$$
\theta_{k+1}=\frac{c \cdot K E_{k}}{\tau}
$$

$$
\omega_{k+1}=\sqrt{\frac{2 C}{I} \cdot K E_{k}}
$$

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## Analysis: Time in two parts


$t_{k 1}$


## Part I

$$
-\tau=I \alpha=I \frac{d \omega}{d t}
$$

for $t=0, \omega=\omega_{k}$

$$
t=\frac{I}{\tau}\left(\omega_{k}-\omega\right)
$$

for $\omega=0, t=t_{k 1}$

$$
t_{k 1}=\frac{I \omega_{k}}{\tau}
$$

## Part 2

$$
\tau=I \alpha=I \frac{d \omega}{d t}
$$

$$
\text { For } t=0, \omega=0
$$

$$
t=\frac{I}{\tau} \omega
$$

For $\omega=\omega_{k}, t=t_{k 2}$

$$
t_{k 2}=\frac{I \omega_{k}}{\tau}
$$

## Time between successive collisions is:

$$
\begin{aligned}
t_{k}=t_{k 1}+t_{k 2} & =2 \frac{I w_{k}}{\tau} \\
& =\frac{2 I}{\tau} \sqrt{\frac{2 c \cdot K E_{k-1}}{I}} \quad \text { where } \quad K E_{k-1}=\frac{1}{2} I \omega_{k-1}^{2} \\
& =\frac{2 I}{\tau} \sqrt{c w_{k-1}^{2}} \\
& =\sqrt{c} t_{k-1} \\
& =(\sqrt{c})^{k-1} t_{0} \quad \text { since } \quad \frac{t_{k}}{t_{k-1}}=\sqrt{c}
\end{aligned}
$$

## Set up and obtaining Spectrograms



## Expectations for times between Successive Amplitudes

- We preliminarily observed collision time periods to be geometric progressions, or constant consecutive ratios of terms.
- We obtain the expression for the time period between collisions:

$$
t_{k}=(\sqrt{c})^{k-1} t_{0}
$$

- We have investigated the time period over the amplitude since it is easier to quantify with our available resources.



## Nature of Collisions

We model each collision sound pulse at time $\boldsymbol{t}$ as a delta function of the form $\boldsymbol{\delta}$ ( $x-x_{t}$ )


When decomposed into its fourier components, the delta function produces a large range of frequencies which we see in the figure.


## Variation of fraction of Energy loss ( $\mathrm{e}^{2}$ )

- $e=$ coefficient of restitution
- $S_{y}=$ dynamic yield strength (dynamic "elastic limit")
- $E^{\prime}=$ effective elastic modulus
- $\rho=$ density
- $v=$ velocity at impact
- $\mu=$ Poisson's ratio
$e=3.1\left(\frac{S_{\mathrm{y}}}{1}\right)^{\frac{5}{8}}\left(\frac{1}{E^{\prime}}\right)^{\frac{1}{2}}\left(\frac{1}{v}\right)^{\frac{1}{4}}\left(\frac{1}{\rho}\right)^{\frac{1}{8}}$

Notice that decrease in velocity leads to increase in $\mathbf{e}$ which is a decrease in the fraction of energy lost.

This leads to the following conclusion:
Locally, (i.e for small number of consecutive collisions), $\boldsymbol{e}$ is constant and our simplified model works for the time analysis. But over time e changes enough to produce a noticeable change

## Display of local character in Data




Time (not to scale)

## Display of local character in Data




## Further Data with similar character



## Original



## Calculation of Exact behaviour

By incorporating $e=3.1\left(\frac{S_{\mathrm{y}}}{1}\right)^{\frac{5}{8}}\left(\frac{1}{E^{\prime}}\right)^{\frac{1}{2}}\left(\frac{1}{v}\right)^{\frac{1}{4}}\left(\frac{1}{\rho}\right)^{\frac{1}{8}}$
into our original equations instead of using constant $\mathbf{e}$ we can calculate the exact behaviour of the frequency. This can be easily modelled numerically.

## Summary of Assumptions

- Constant torque



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- Rigid body (no change in shape)


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- Constant torque
- Rigid body (no change in shape)
- Small angle approximation has been made freely in the analysis


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## Further aspects of consideration

- Effect of temperature.
- Effect of modulus of material.
- Effect of liquid in glasses.


## Thank you.

## End trail (Hard to Reproduce)



